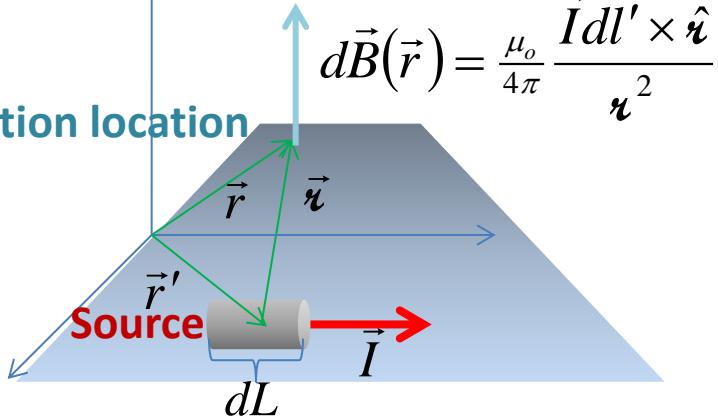
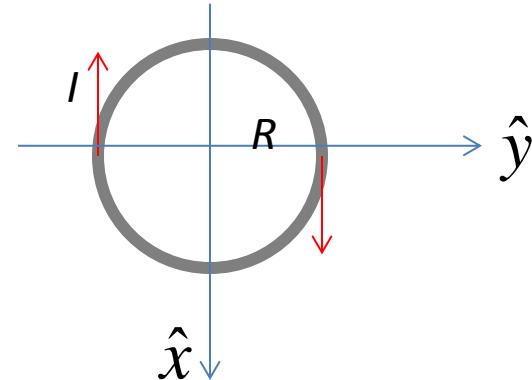


Mon.	(C 17) 5.2 Biot-Savart Law T5 Quiver Plots	
Tues.		HW7
Wed.	(C 21.6-7,.9) 1.3.4-1.3.5, 1.5.2-1.5.3, 5.3.1-.3.2 Div & Curl B	
Fri.	(C 21.6-7,.9) 5.3.3-.3.4 Applications of Ampere's Law	
Mon.	1.6, 5.4.1-.4.2 Magnetic Vector Potential	
Wed.	5.4.3 Multipole Expansion of the Vector Potential	
Thurs.		HW8
Fri.	Review	

Biot-Savart Law

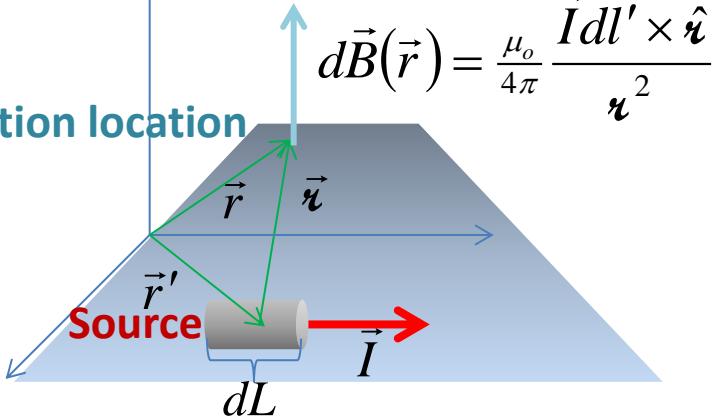


Exercise: field in center of clockwise circle of current, of radius R, lying in the x-y plane.

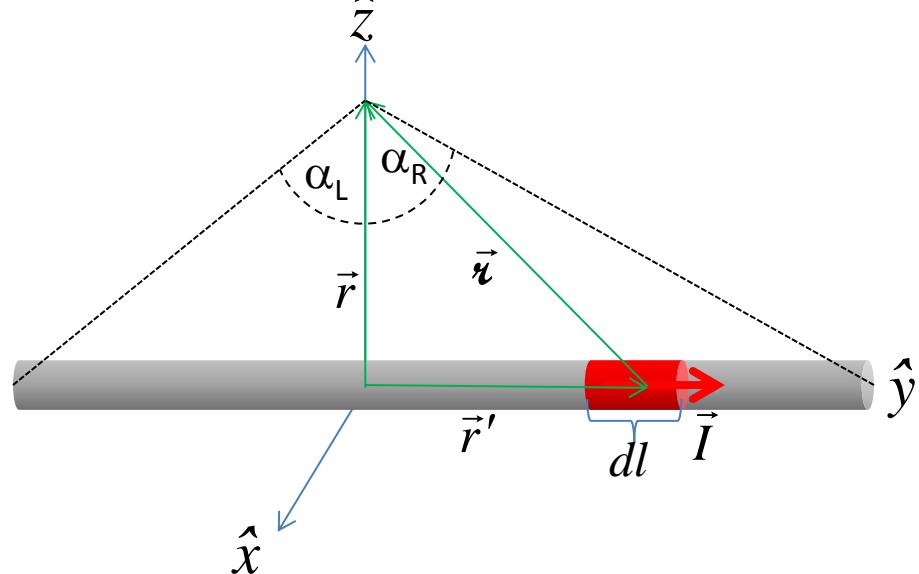


Exercise: what's the contribution of just a third of the circle?

Biot-Savart Law



Example: field due to a line segment of current
(in terms of angles and $|r|$)



Biot-Savart Law

$$\vec{B}(\vec{r}) = \frac{\mu_0}{4\pi} I \int \frac{d\vec{\ell}' \times \hat{\vec{r}}}{r^2}$$

$$d\vec{\ell}' \times \hat{\vec{r}} = |d\vec{\ell}' \times \hat{\vec{r}}| \hat{x}$$

$$|d\vec{\ell}' \times \hat{\vec{r}}| = |d\ell| |\hat{\vec{r}}| \sin(\alpha + 90^\circ) = d\ell \cos \alpha$$

$$\vec{B}(\vec{r}) = \frac{\mu_0}{4\pi} I \int \frac{d\ell \cos \alpha}{r^2} \hat{x}$$

$$\frac{d\ell}{r^2} = \frac{1}{r} \frac{r^2 d(\ell/r)}{r^2} = \frac{1}{r} \cdot \left(\frac{r}{r} \right)^2 \cdot d\left(\frac{\ell}{r} \right) = \frac{1}{r} \cdot \left(\frac{r}{r} \right)^2 \cdot d\left(\frac{r'}{r} \right)$$

$(\cos \alpha)^2 \ d(\tan \alpha)$

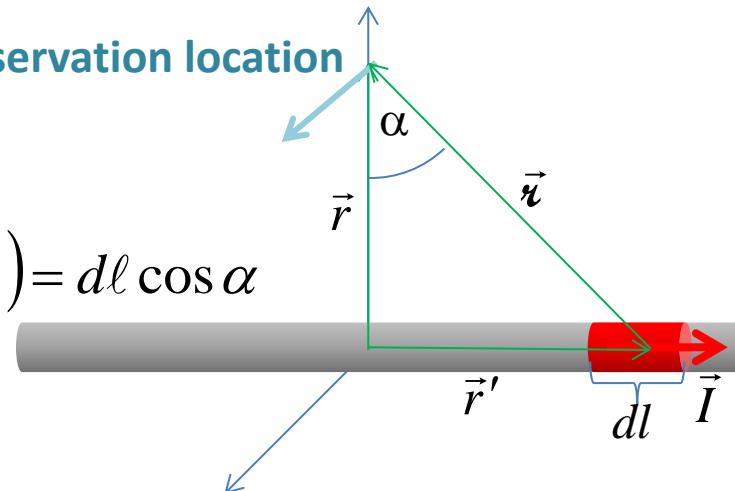
$$d(\tan \alpha) = \frac{d\alpha}{\cos^2 \alpha}$$

$$\begin{aligned} \vec{B}(\vec{r}) &= \frac{\mu_0}{4\pi} I \int \frac{1}{r} (\cos \alpha)^2 \frac{d\alpha}{(\cos \alpha)^2} \cos \alpha \hat{x} = \frac{\mu_0}{4\pi} I \int \frac{(\cos \alpha) d\alpha}{r} \hat{x} = \frac{\mu_0}{4\pi} \frac{I}{r} \int_{\alpha_R}^{\alpha_L} d(\sin \alpha) \hat{x} \\ &= \frac{\mu_0}{4\pi} \frac{2I}{r} (\sin(\alpha_L) - \sin(\alpha_R)) \hat{x} \end{aligned}$$

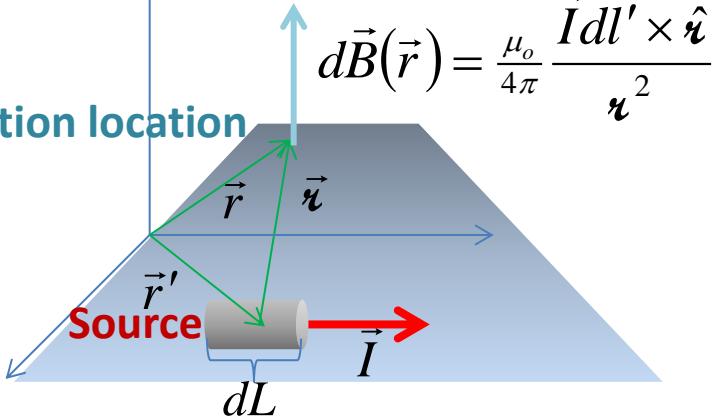
Field of Finite Wire

(Book's approach)

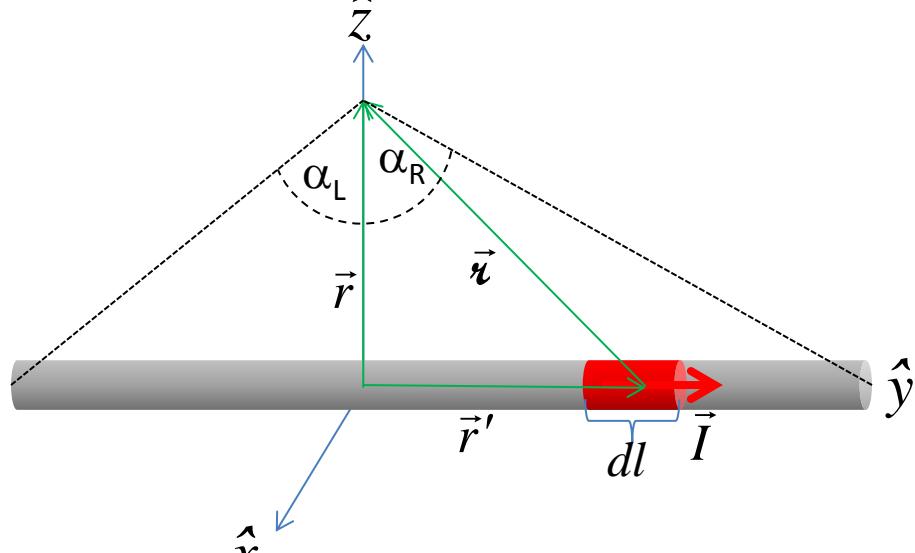
Observation location



Biot-Savart Law

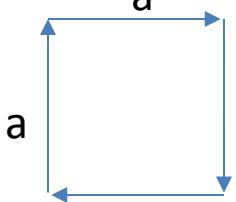


Example: field due to a line segment of current (in terms of angles and $|r|$)

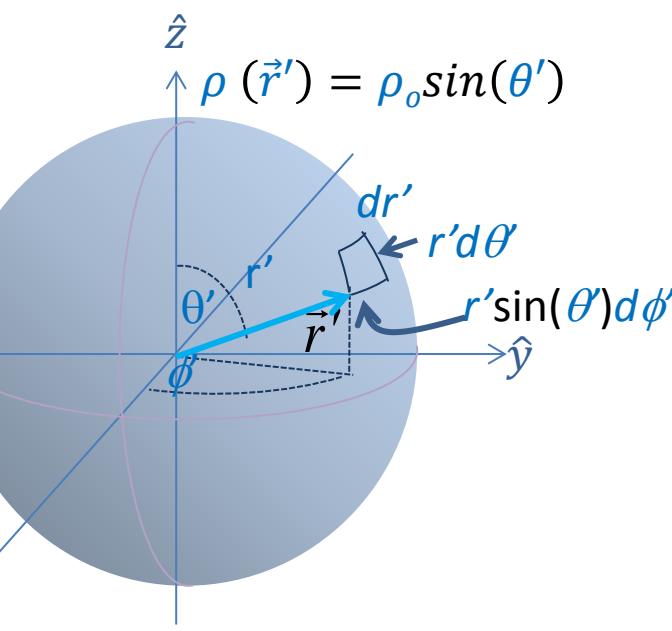


$$\vec{B}(\vec{r}) = \frac{\mu_0}{4\pi} \frac{2I}{r} (\sin(\alpha_L) - \sin(\alpha_R)) \hat{x}$$

Exercise: field at center of square, of side a , carrying current I



What's the magnetic field in the center of a sphere with charge density ρ rotating about z at ω .



$$d\vec{B}(\vec{r}) = \frac{\mu_0}{4\pi} \frac{\vec{I}dl' \times \hat{r}}{r'^2} = \frac{\mu_0}{4\pi} \frac{\vec{I}dl' \times \vec{r}}{r'^3}$$

so, we want to figure out exactly what $d\vec{I}dl$ is, what \vec{r} is, and what their cross product is, then what r is.

One way of approaching $d\vec{I}dl$ is that it's really $\vec{v}dq$, since that's rotating at constant rate $\vec{v} = \omega(r' \sin \theta)\hat{\phi}$ and the bit of charge is $dq = \rho d\tau'$. In our case,

$$dq = (\rho_0 \sin \theta')(r'^2 dr' \sin \theta' d\theta' d\phi')$$

So putting together $d\vec{I}dl = \vec{v}dq = \vec{v}\rho d\tau'$ gives us

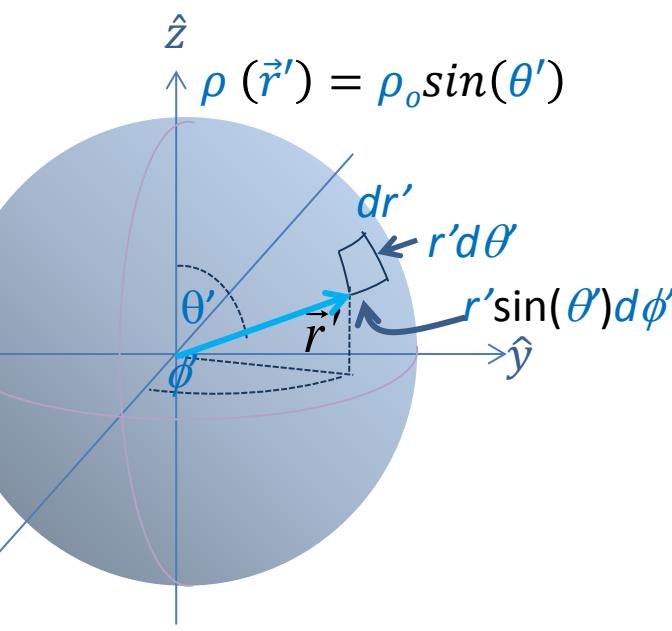
$$d\vec{I}dl = (\omega r' \sin \theta' \hat{\phi})(\rho_0 \sin \theta')(r'^2 \sin \theta' dr' d\theta' d\phi')$$

$$\vec{r} = \cancel{\langle x - r' \cos \phi' \sin \theta', y - r' \sin \phi' \sin \theta', z - r' \cos \theta' \rangle}$$

$$|\vec{r}| = r'$$

For the sake of doing the cross product, we'll want to say $\hat{\phi} = \langle -\sin \phi', \cos \phi', 0 \rangle$

What's the magnetic field in the center of a sphere with charge density ρ rotating about z at ω .



$$d\vec{B}(\vec{r}) = \frac{\mu_0}{4\pi} \frac{\vec{I}dl' \times \hat{\vec{r}}}{r'^2} = \frac{\mu_0}{4\pi} \frac{\vec{I}dl' \times \hat{\vec{r}}}{r'^3}$$

$$d\vec{I}dl = \rho \omega r'^3 \sin^3 \theta' d\theta' dr' d\phi' \hat{\vec{\phi}}$$

$$\hat{\vec{r}} = -r' \langle \cos \phi' \sin \theta', \sin \phi' \sin \theta', \cos \theta' \rangle$$

$$|\hat{\vec{r}}| = r'$$

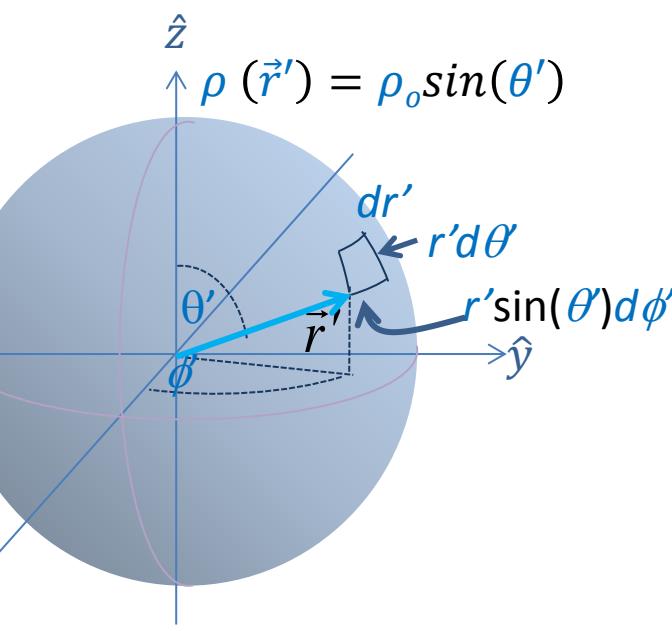
$$\hat{\vec{\phi}} = \langle -\sin \phi', \cos \phi', 0 \rangle$$

$$\hat{\vec{\phi}} \times \hat{\vec{r}} = -r' \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ -\sin \phi' & \cos \phi' & 0 \\ \cos \phi' \sin \theta' & \sin \phi' \sin \theta' & \cos \theta' \end{vmatrix}$$

$$\hat{\vec{\phi}} \times \hat{\vec{r}} = -r' (\hat{x} \cos \phi' (\cos \theta') + \hat{y} \sin \phi' (\cos \theta') + \hat{z} (-\sin \phi' \sin \theta') - \cos \phi' (\cos \phi' \sin \theta'))$$

$$\hat{\vec{\phi}} \times \hat{\vec{r}} = -r' (\hat{x} \cos \phi' (\cos \theta') + \hat{y} \sin \phi' (\cos \theta') - \hat{z} (\sin \theta'))$$

What's the magnetic field in the center of a sphere with charge density ρ rotating about z at ω .



$$d\vec{B}(\vec{r}) = \frac{\mu_0}{4\pi} \frac{\vec{I} dl' \times \hat{\vec{r}}}{r'^2} = \frac{\mu_0}{4\pi} \frac{\vec{I} dl' \times \hat{\vec{r}}}{r'^3}$$

$$d\vec{I} dl = \rho_0 \omega r'^3 \sin^3 \theta' d\theta' dr' d\phi' \hat{\phi}$$

$$\hat{\vec{r}} = -r' \langle \cos \phi' \sin \theta', \sin \phi' \sin \theta', \cos \theta' \rangle$$

$$|\hat{\vec{r}}| = r'$$

$$\hat{\phi} = \langle -\sin \phi', \cos \phi', 0 \rangle$$

$$\hat{\phi} \times \hat{\vec{r}} = -r' (\hat{x} \cos \phi' \cos \theta' + \hat{y} \sin \phi' \cos \theta' - \hat{z} \sin \theta')$$

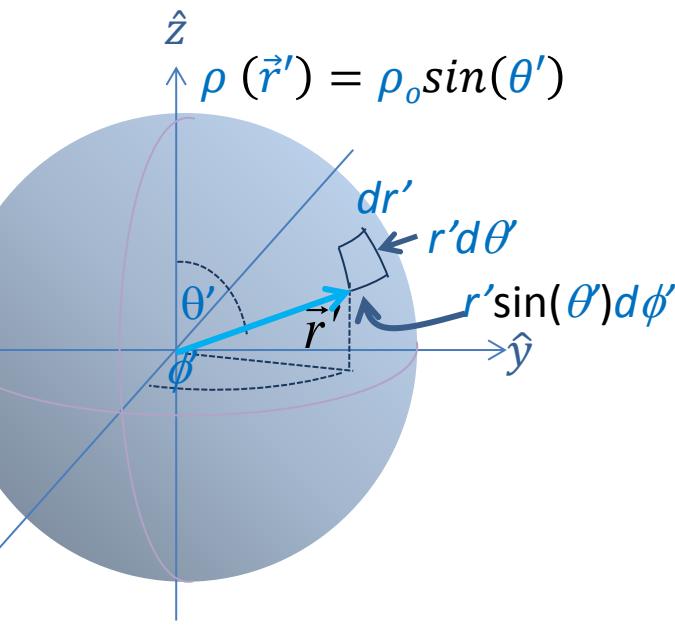
$$\vec{B}(0) = \frac{\mu_0}{4\pi} \int \frac{(\rho_0 \omega r'^3 \sin^3 \theta' d\theta' dr' d\phi') (-r') (\hat{x} \cos \phi' \cos \theta' + \hat{y} \sin \phi' \cos \theta' - \hat{z} \sin \theta')}{r'^3}$$

$$\vec{B}(0) = -\frac{\mu_0}{4\pi} \rho_0 \omega \int (r' \sin^3 \theta' d\theta' dr' d\phi') (\cos \phi' \cos \theta' \hat{x} + \sin \phi' \cos \theta' \hat{y} - \sin \theta' \hat{z})$$

Symmetry tells us there will only be a z-component in the end, so skipping the other two:

$$\vec{B}(0) = -\frac{\mu_0}{4\pi} \rho_0 \omega \int (r' \sin^3 \theta' d\theta' dr' d\phi') (-\sin \theta') \hat{z}$$

What's the magnetic field in the center of a sphere with charge density ρ rotating about z at ω .



$$\vec{B}(0) = \frac{\mu_o}{4\pi} \rho_o \omega \int (r' \sin^4 \theta' d\theta' dr' d\phi') \hat{z}$$

$$\vec{B}(0) = \frac{\mu_o}{4\pi} \rho_o \omega \int_0^{r'=R} \int_0^{\theta=\pi} \int_0^{\phi=2\pi} r' \sin^4 \theta' d\theta' dr' d\phi' \hat{z}$$

$$\vec{B}(0) = \frac{\mu_o}{4\pi} \rho_o \omega (2\pi) \int_0^{r'=R} \int_0^{\theta=\pi} r' \sin^4 \theta' d\theta' dr' \hat{z}$$

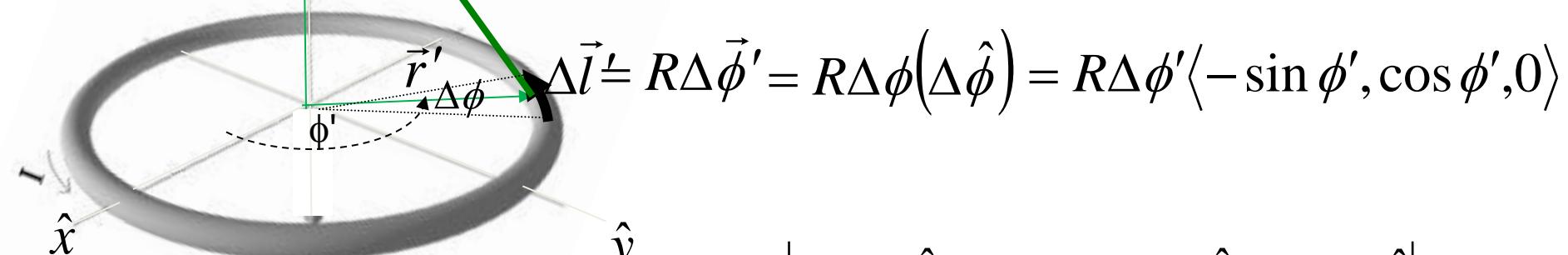
$$\vec{B}(0) = \frac{\mu_o}{2} \rho_o \omega \left(\frac{1}{2} R^2 \right) \int_0^{\theta=\pi} \sin^4 \theta' d\theta' \hat{z}$$

$$\vec{B}(0) = \frac{\mu_o}{4} \rho_o \omega R^2 \left(\frac{3\pi}{8} \right) \hat{z} = \frac{\mu_o 3\pi R^2 \rho_o \omega}{64} \hat{z}$$

B field of loop – on axis $\vec{r} = \langle 0, 0, z \rangle$

$$\Delta \vec{B} = \frac{\mu_0}{4\pi} \frac{I \Delta \vec{l}' \times \hat{\vec{r}}}{r^2} = \frac{\mu_0}{4\pi} \frac{I \Delta \vec{l}' \times \vec{r}}{r^3}$$

$$\vec{r} = \vec{r} - \vec{r}' = \langle -R \cos \theta', -R \sin \theta', z \rangle$$

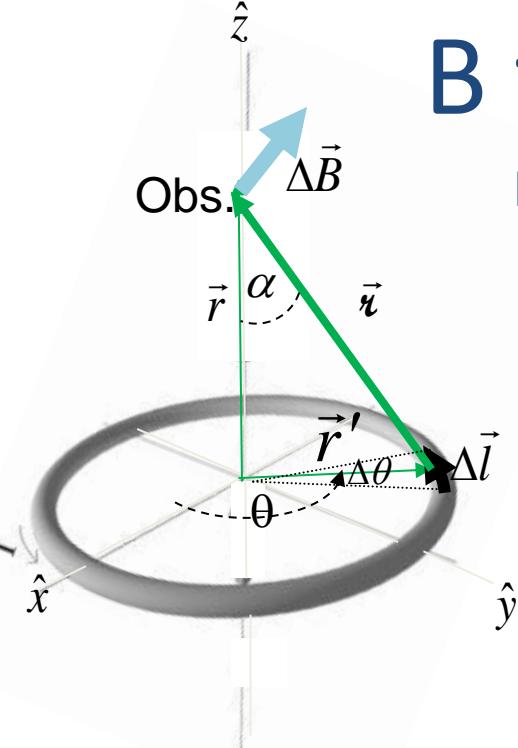


$$\Delta \vec{l}' \times \vec{r} = \det \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ -R \Delta \phi' \sin \phi' & R \Delta \phi' \cos \phi' & 0 \\ -R \cos \phi' & -R \sin \phi' & z \end{vmatrix}$$

$$\Delta \vec{l}' \times \vec{r} = R \Delta \phi' \langle z \cos \phi', z \sin \phi', R \rangle$$

B field of loop – on axis

Obs.



By Symmetry (could do math, but why bother)

$$B_x = \frac{\mu_0}{4\pi} IRz \int_{\theta'=0}^{\theta'=2\pi} \frac{\cos \theta' d\theta'}{(r^2 + R^2)^{\frac{3}{2}}} = \dots = 0$$

$$B_y = \frac{\mu_0}{4\pi} IRz \int_{\theta'=0}^{\theta'=2\pi} \frac{\sin \theta' d\theta'}{(r^2 + R^2)^{\frac{3}{2}}} = \dots = 0$$

$$B_z = \frac{\mu_0}{4\pi} IR^2 \int_{\theta'=0}^{\theta'=2\pi} \frac{d\theta'}{(r^2 + R^2)^{\frac{3}{2}}} = \frac{\mu_0}{4\pi} \frac{IR^2 z}{(z^2 + R^2)^{\frac{3}{2}}} \int_{\theta=0}^{\theta=2\pi} d\theta$$

Rephrase in terms of angle α ...

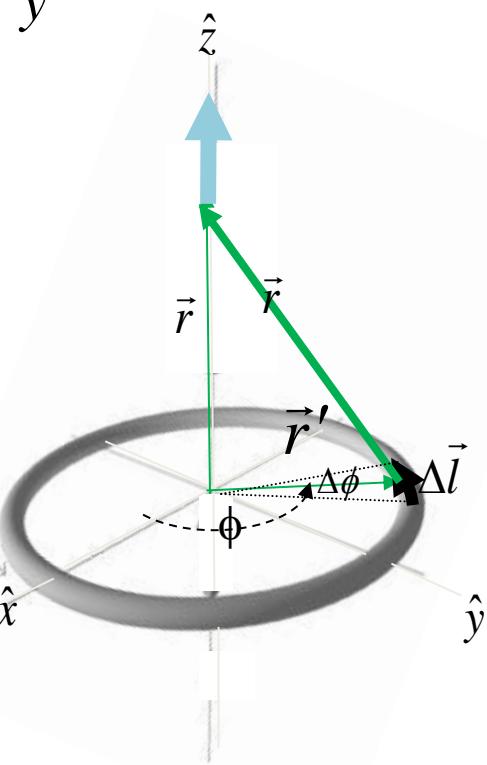
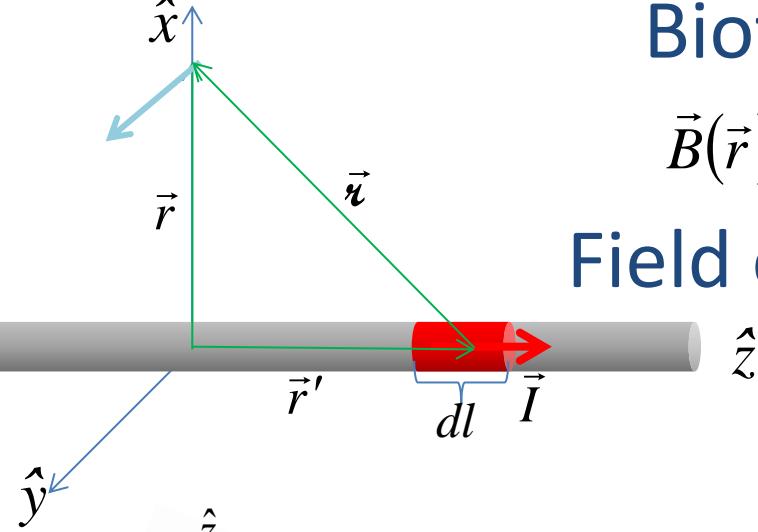
$$= \frac{\mu_0}{4\pi} \frac{IR^2}{(z^2 + R^2)^{\frac{3}{2}}} 2\pi$$

In special case of $z = 0$, reduces to...

Biot-Savart Law

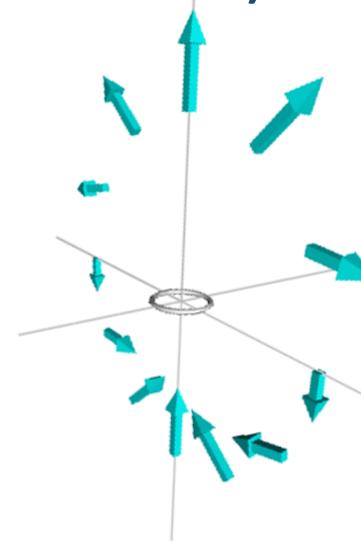
$$\vec{B}(\vec{r}) = \frac{\mu_0}{4\pi} I \int \frac{d\vec{l}' \times \hat{r}}{r^2}$$

Field of Infinite Wire



of a Loop (on axis)

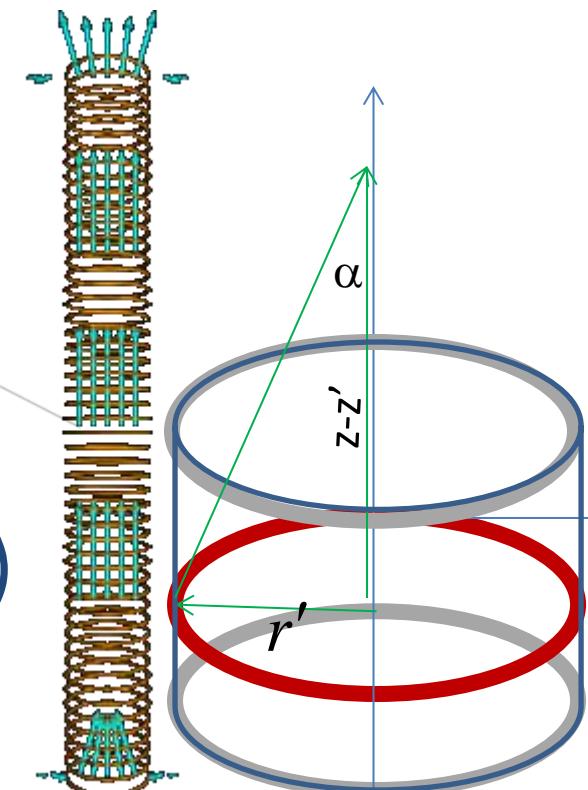
$$\frac{\mu_0}{2} \frac{IR^2}{(z^2 + R^2)^{\frac{3}{2}}} \hat{z}$$



of a Solenoid (on axis)

$$\frac{\mu_0}{2} \frac{IN}{L} (\cos \alpha_{bottom} - \cos \alpha_{top}) \hat{z}$$

Cartesian	Polar
$\frac{\mu_0}{4\pi} \frac{2I}{x} \hat{y} = \frac{\mu_0}{4\pi} \frac{2I}{s} \hat{\phi}$	$\frac{\mu_0}{4\pi} \frac{2I}{s} \hat{\phi}$



Biot-Savart Law of a disc of nested Loops (on axis)

$$\vec{B}(z)_{nest} = \sum_{loops} \vec{B}(z)_{loop,i}$$

$$\vec{B}(z)_{nest} = \sum_{loops} \frac{\mu_o}{2} \frac{Ir'^2}{(z^2 + r'^2)^{\frac{3}{2}}} \hat{z}$$

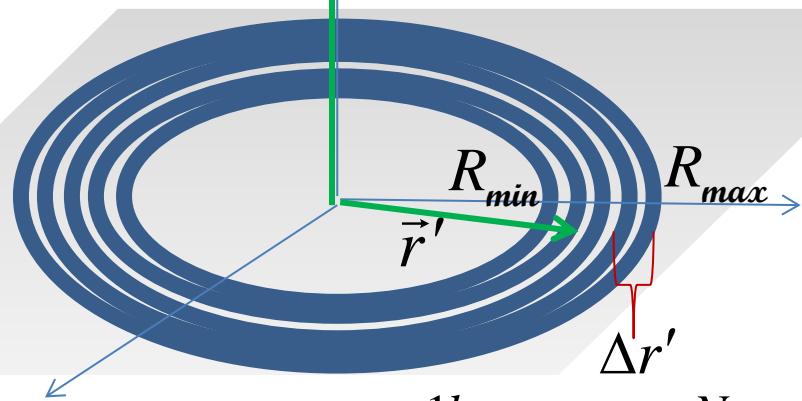
$$\vec{B}(z)_{nest} = \sum_{r'} \frac{\mu_o}{2} \frac{Ir'^2}{(z^2 + r'^2)^{\frac{3}{2}}} \left(\Delta r' \left(\frac{N}{R_{max} - R_{min}} \right) \right) \hat{z}$$

$$\vec{B}(z)_{nest} = \frac{\mu_o}{2} I \frac{N}{R_{max} - R_{min}} \int_{r'=R_{min}}^{R_{max}} \frac{r'^2 dr'}{(z^2 + r'^2)^{\frac{3}{2}}} \hat{z}$$

$$\vec{B}(z)_{nest} = \frac{\mu_o}{2} I \frac{N}{R_{max} - R_{min}} \left[\ln \left(\sqrt{z^2 + r'^2} + r' \right) - \frac{r'}{\sqrt{z^2 + r'^2}} \right]_{R_{min}}^{R_{max}} \hat{z}$$

$$\vec{B}(z)_{nest} = \frac{\mu_o}{2} I \frac{N}{R_{max} - R_{min}} \left[\ln \left(\frac{\sqrt{\left(\frac{z}{R_{max}} \right)^2 + 1} + 1}{\sqrt{\left(\frac{z}{R_{min}} \right)^2 + 1} + 1} \right) + \frac{1}{\sqrt{\left(\frac{z}{R_{min}} \right)^2 + 1}} - \frac{1}{\sqrt{\left(\frac{z}{R_{max}} \right)^2 + 1}} \right] \hat{z}$$

$$\vec{B}(z)_{loop} = \frac{\mu_o}{2} \frac{Ir'^2}{(z^2 + r'^2)^{\frac{3}{2}}} \hat{z}$$



$$\frac{1}{\Delta r'} = \frac{N}{R_{max} - R_{min}}$$

$$1 \text{ loop} = \frac{N}{R_{max} - R_{min}} \Delta r'$$

Biot-Savart Law of a disc of nested Loops (on axis)

$$\vec{B}(z)_{nest} = \sum_{loops} \vec{B}(z)_{loop,i}$$

$$\vec{B}(z)_{nest} = \sum_{loops} \frac{\mu_0}{2} \frac{Ir'^2}{(z^2 + r'^2)^{\frac{3}{2}}} \hat{z}$$

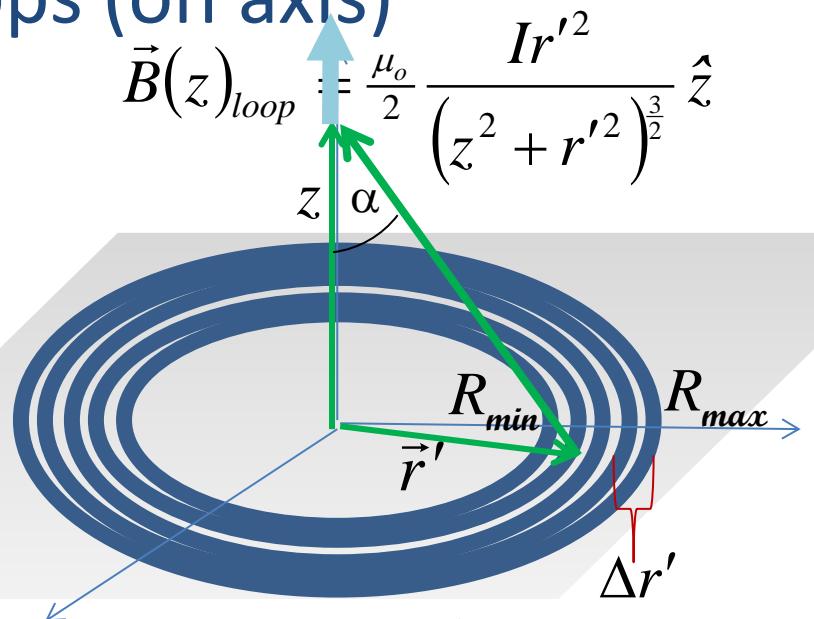
$$\vec{B}(z)_{nest} = \sum_{r'} \frac{\mu_0}{2} \frac{Ir'^2}{(z^2 + r'^2)^{\frac{3}{2}}} \left(\Delta r' \left(\frac{N}{R_{max} - R_{min}} \right) \right) \hat{z}$$

$$\vec{B}(z)_{nest} = \frac{\mu_0}{2} I \int_{r'=R_{min}}^{R_{max}} \frac{r'^2 dr'}{(z^2 + r'^2)^{\frac{3}{2}}} \hat{z}$$

Exercise: translate into integral of trig function of α .

$$\frac{r'^2 dr'}{(z^2 + r'^2)^{\frac{3}{2}}} = \left(\frac{r'}{\sqrt{z^2 + r'^2}} \right)^2 \frac{dr'}{\sqrt{z^2 + r'^2}} = (\sin \alpha)^2 \frac{d(z \tan \alpha)}{(z / \cos \alpha)} = (\sin \alpha)^2 \frac{\cos \alpha d\alpha}{(\cos \alpha)^2}$$

$$= (1 - \cos^2 \alpha) \frac{\cos \alpha d\alpha}{(\cos \alpha)^2} = \left(\frac{d\alpha}{\cos \alpha} - \cos \alpha d\alpha \right)$$



$$\frac{1loop}{\Delta r'} = \frac{N}{R_{max} - R_{min}}$$

$$1loop = \frac{N}{R_{max} - R_{min}} \Delta r'$$

$$= (1 - \cos^2 \alpha) \frac{\cos \alpha d\alpha}{(\cos \alpha)^2}$$

Exercise:

disc of uniform surface charge density, K , spinning at ω

$$\vec{B}(z)_{nest} = \sum_{loops} \vec{B}(z)_{loop,i}$$

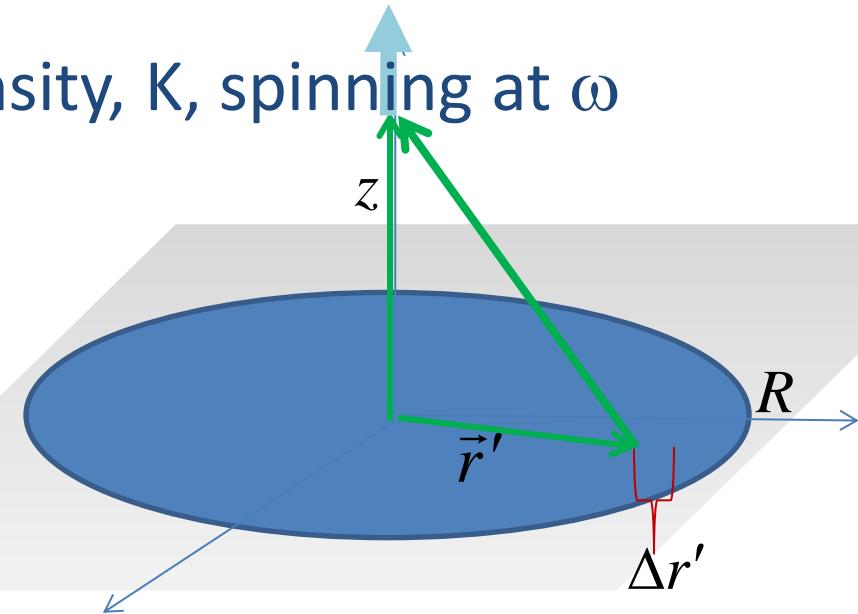
$$\vec{B}(z)_{nest} = \sum_{loops} \frac{\mu_0}{2} \frac{(dI)r'^2}{(z^2 + r'^2)^{\frac{3}{2}}} \hat{z}$$

Recall: $K = \frac{dI}{dl_{\perp}} \Rightarrow dI = K dl_{\perp}$

Where dl_{\perp} is a small step perpendicular to current flow, i.e. dr'

From the homework that you've just done, $K = \sigma v = \sigma \omega r'$ so, $dI = \sigma \omega r' dr'$

$$\begin{aligned} \vec{B}(z)_{nest} &= \int_0^R \frac{\mu_0}{2} \frac{(\sigma \omega r' dr') r'^2}{(z^2 + r'^2)^{\frac{3}{2}}} \hat{z} = \frac{\mu_0}{2} \sigma \omega \int_0^R \frac{r'^3 dr'}{(z^2 + r'^2)^{\frac{3}{2}}} \hat{z} = \frac{\mu_0}{2} \sigma \omega z \int_0^{\frac{R}{z}} \frac{\left(\frac{r'}{z}\right)^3 d\left(\frac{r'}{z}\right)}{\left(1 + \left(\frac{r'}{z}\right)^2\right)^{\frac{3}{2}}} \hat{z} \\ &= \frac{\mu_0}{4} \sigma \omega z \int_0^{\frac{R}{z}} \frac{\left(\frac{r'}{z}\right)^2 d\left(\frac{r'}{z}\right)^2}{\left(1 + \left(\frac{r'}{z}\right)^2\right)^{\frac{3}{2}}} \hat{z} = \frac{\mu_0}{4} \sigma \omega z \int_0^{\left(\frac{R}{z}\right)^2} \frac{\xi d\xi}{(1 + \xi)^{\frac{3}{2}}} \hat{z} = \frac{\mu_0}{4} \sigma \omega z \left[\frac{2(\xi + 2)}{(1 + \xi)^{\frac{1}{2}}} \right]_0^{\left(\frac{R}{z}\right)^2} \\ &= \frac{\mu_0}{2} \sigma \omega z \left[\frac{\left(\frac{R}{z}\right)^2 + 2}{\left(1 + \left(\frac{R}{z}\right)^2\right)^{\frac{1}{2}}} - 2 \right] \end{aligned}$$

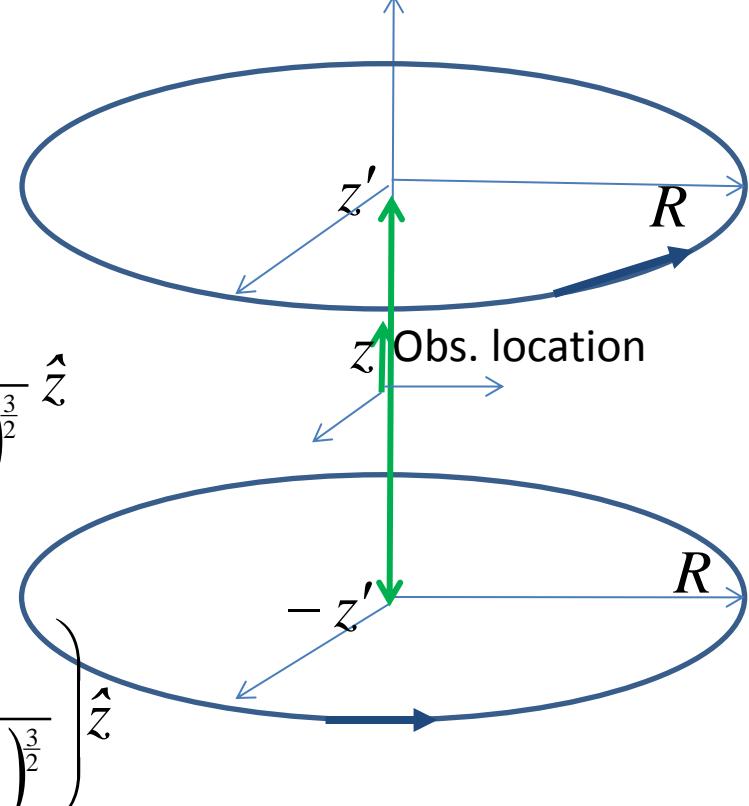


Helmholtz Coils

$$\vec{B}(z)_{Hc} = \vec{B}(z)_{loop.\text{top}} + \vec{B}(z)_{loop.\text{bottom}}$$

$$\vec{B}(z)_{Hc} = \frac{\mu_o}{2} \frac{IR^2}{((z - z')^2 + R^2)^{\frac{3}{2}}} \hat{z} + \frac{\mu_o}{2} \frac{IR^2}{((z + z')^2 + R^2)^{\frac{3}{2}}} \hat{z}$$

$$\vec{B}(z)_{Hc} = \frac{\mu_o}{2} IR^2 \left(\frac{1}{((z - z')^2 + R^2)^{\frac{3}{2}}} + \frac{1}{((z + z')^2 + R^2)^{\frac{3}{2}}} \right) \hat{z}$$



Why they're useful

$$\left. \frac{dB(z)_{Hc}}{dz} \right|_{z=0} = -\frac{3\mu_o}{2} IR^2 \left(\frac{(0 - z')}{((0 - z')^2 + R^2)^{\frac{5}{2}}} + \frac{(0 + z')}{((0 + z')^2 + R^2)^{\frac{5}{2}}} \right) = 0$$

and

$$\left. \frac{d^2 B(z)_{Hc}}{dz^2} \right|_{z=0} = 0 \quad \text{if } z' = \frac{R}{2} \quad \text{Produces fairly uniform field near middle}$$

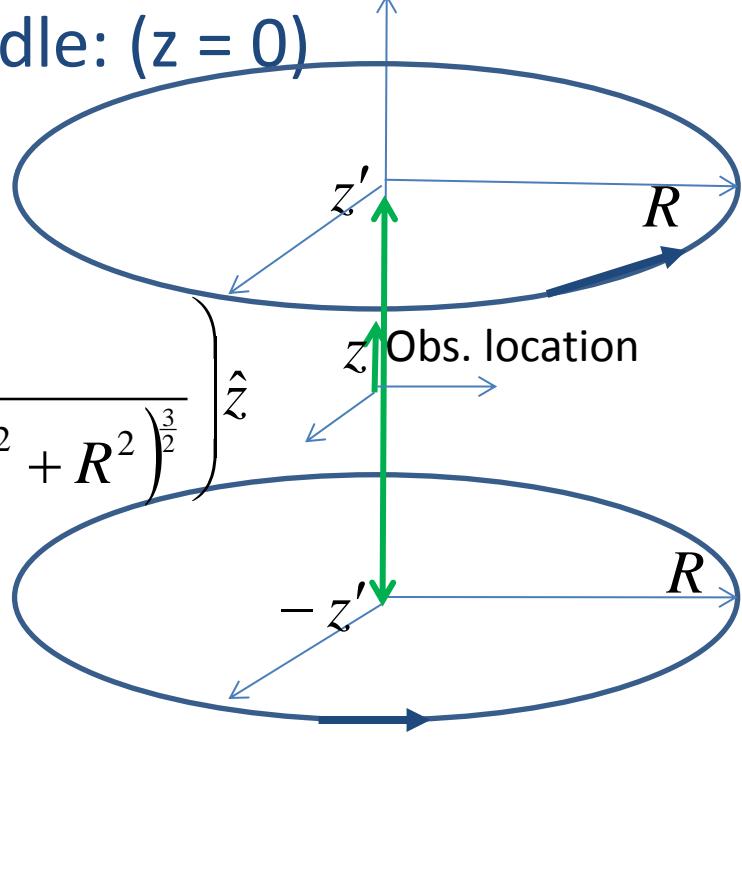
Helmholtz Coils If $z' = R/2$ field in the middle: ($z = 0$)

$$\vec{B}(0)_{Hc} = \vec{B}(0)_{loop.top} + \vec{B}(0)_{loop.bottom}$$

$$\vec{B}(0)_{Hc} = \frac{\mu_o}{2} IR^2 \left(\frac{1}{\left((0 - R/2)^2 + R^2 \right)^{\frac{3}{2}}} + \frac{1}{\left((0 - R/2)^2 + R^2 \right)^{\frac{3}{2}}} \right) \hat{z}$$

$$\vec{B}(0)_{Hc} = \frac{\mu_o}{2} IR^2 \left(\frac{2}{\left(5R^2 / 4 \right)^{\frac{3}{2}}} \right) \hat{z}$$

$$\vec{B}(0)_{Hc} = \left(\frac{4}{5} \right)^{\frac{3}{2}} \frac{\mu_o I}{R} \hat{z}$$



Mon.	(C 17) 5.2 Biot-Savart Law T5 Quiver Plots	
Tues.		HW7
Wed.	(C 21.6-7,.9) 1.3.4-1.3.5, 1.5.2-1.5.3, 5.3.1-.3.2 Div & Curl B	
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Fri.	Review	