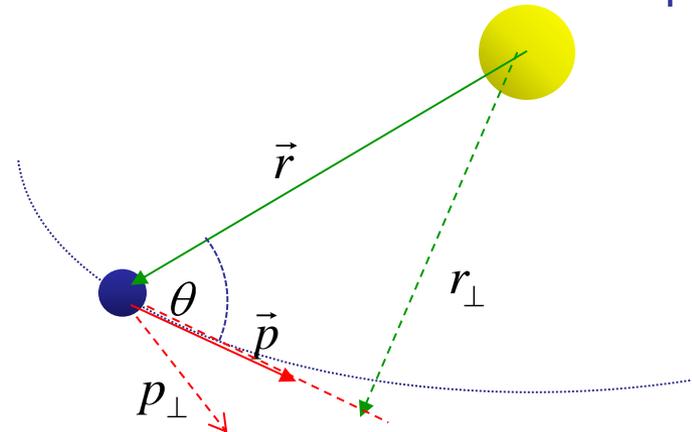


Fri.	11.10 Quantization, Quiz 11 Lecture Evals	RE 11.e
Mon.	Review for Final (1-11)	HW11: Pr's 39, 57, 64, 74, 78
Sat.	9 a.m.	Final Exam (Ch. 1-11)

# Using Angular Momentum

The measure of motion *about* a point



## Magnitude and Direction

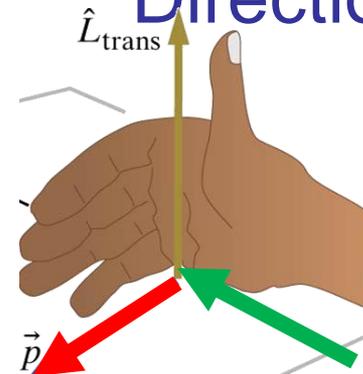
$\vec{L} =$

$\langle \quad \quad \quad \rangle$

Magnitude

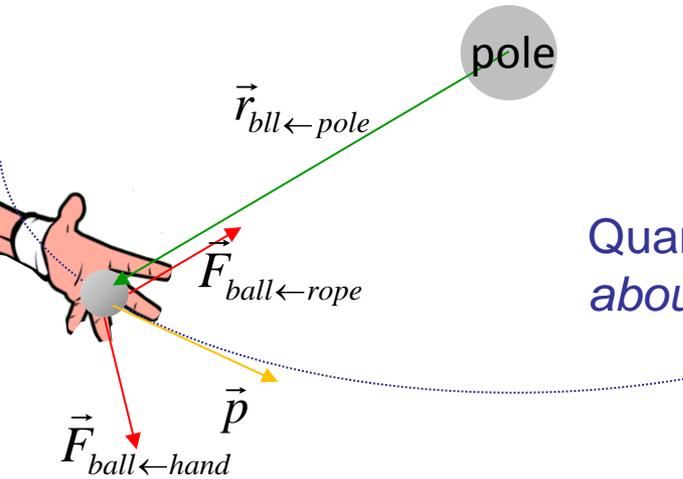
$$|L| = |p_{\perp}|r = |p|r_{\perp} = |p|r \sin(\theta)$$

Direction



Orient Right hand so fingers curl from the axis and with motion, then thumb points in direction of angular momentum.

# Angular Momentum Principle



$$\frac{d}{dt} \vec{L}_{(about)A} = \sum_{net} \vec{\tau}_{(about)A}$$

Quantifies motion  
*about* a point

**Torque**

where,  $\vec{\tau}_{(about)A} \equiv \vec{r}_{(from)A} \times \vec{F}$

Interaction that  
changes motion  
*about* a point

**Magnitude**

(yet another cross product)

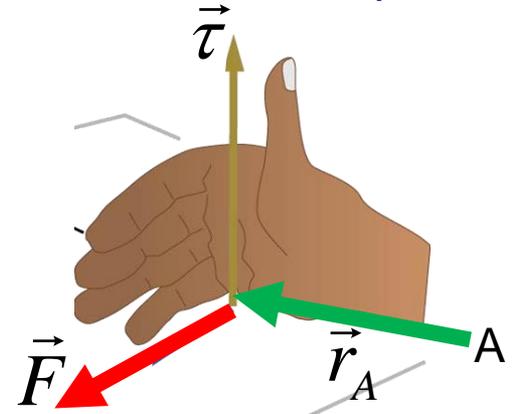
$$|\tau_A| = |r_A| |F_{\perp}| = |r_{A\perp}| |F|$$

$$|\tau_A| = |r_A| |F| \sin \theta = (|r_A| \sin \theta) |F| = |r_{A\perp}| |F| \sin \theta$$

Making sense of the  
factors and cross-product

**Direction**

(yet another cross product)



# Zero-Torque Systems

## Spinning Skater



Initial



Final

$$\vec{L}_{rot.f} - \vec{L}_{rot.i} = \vec{\tau}_{ave} \Delta t$$

$$\vec{L}_{rot.f} - \vec{L}_{rot.i} \approx 0$$

$$\vec{L}_{rot.i} \approx \vec{L}_{rot.f}$$

mass farther from axis:  $I_i \vec{\omega}_i \approx I_f \vec{\omega}_f$  mass closer to axis:  
 $I_i$  larger  $I_f$  smaller  
 $\omega_i$  smaller  $\omega_f$  larger

# Three Fundamental Principles

Angular Momentum:

$$\frac{d}{dt} \vec{L}_{(about)A} = \sum_{net} \vec{\tau}_{(about)A}$$

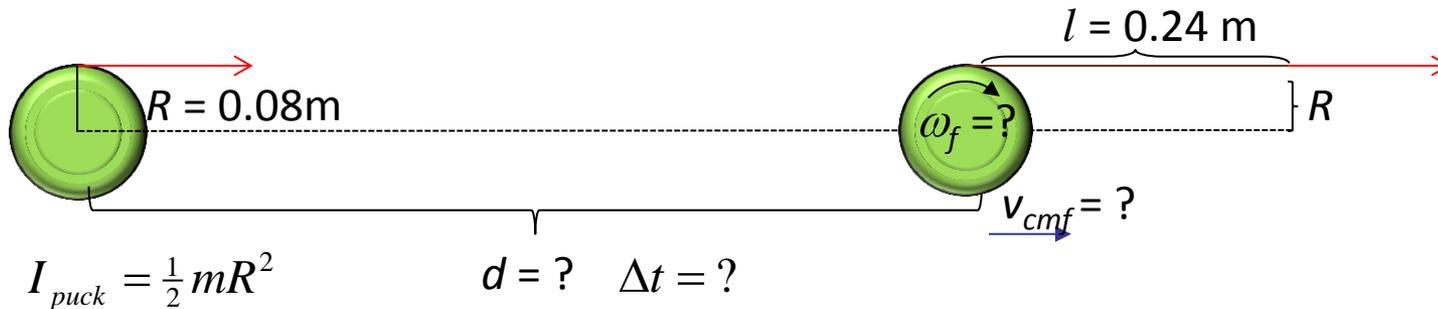
(linear) Momentum:

$$\frac{d}{dt} \vec{p} = \sum_{net} \vec{F}$$

Energy:

$$\Delta E = \sum_{net} W$$

Example all three together! Say we have a uniform 0.4 kg puck with an 8 cm radius. A 24 cm string is initially wrapped around its circumference. If it's on a frictionless surface and a 10 N force is applied to the end of the string until it's unwound...



a. What will be its rate of rotation when the string is fully unwound?

Energy Principle

$$\Delta E_{total} = W$$

$$\omega_f = \sqrt{\frac{2Fl}{I}}$$

$$\Delta K_{trans} + \Delta E_{int} = \vec{F} \cdot \Delta \vec{r}_F$$

$$\omega_f = \sqrt{\frac{2Fl}{\frac{1}{2} R^2 m}} = \frac{2}{R} \sqrt{\frac{Fl}{m}}$$

$$\vec{F} \cdot \Delta \vec{r}_{cm} + \Delta K_{rot} = F(d + l)$$

$$\cancel{Fd} + \left| \frac{1}{2} I \omega_f^2 - \frac{1}{2} I \omega_i^2 \right| = \cancel{Fd} + Fl$$

$$\frac{1}{2} I \omega_f^2 = Fl$$

# Three Fundamental Principles

Angular Momentum:

$$\frac{d}{dt} \vec{L}_{(about)A} = \sum_{net} \vec{\tau}_{(about)A}$$

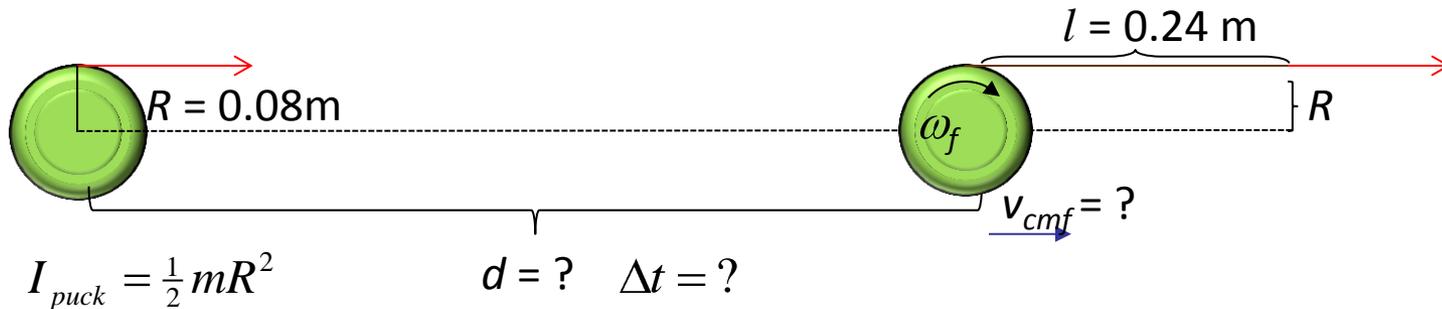
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a. What will be its rate of rotation when the string is fully unwound?  $\omega_f = \frac{2}{R} \sqrt{\frac{Fl}{m}}$

b. How long was the force applied? Angular Momentum Principle (axis through final location of cm)

$$\Delta \vec{L} = \int \vec{\tau} dt$$

$$\vec{L}_f - \vec{L}_i = \vec{\tau} \Delta t$$

$$(\vec{L}_{f,trans} + \vec{L}_{f,rot}) - (\vec{L}_{i,trans} + \vec{L}_{i,rot}) = \vec{\tau} \Delta t$$

$$\Rightarrow \Delta t = \frac{I \omega_f}{RF} = \frac{\frac{1}{2} R^2 m \omega_f}{RF} = \frac{R m \omega_f}{2F} = \frac{R m \frac{2}{R} \sqrt{\frac{Fl}{m}}}{2F} = \sqrt{\frac{ml}{F}}$$

$I \vec{\omega}_f = |\vec{r}_{F-a} \times \vec{F}| \Delta t$  Torque and final angular velocity in  $-z$  direction

$$I \omega_f = |RF| \Delta t \Rightarrow \Delta t = \frac{I \omega_f}{RF} = \frac{\frac{1}{2} R^2 m \omega_f}{RF} = \frac{R m \omega_f}{2F} = \frac{R m \frac{2}{R} \sqrt{\frac{Fl}{m}}}{2F} = \sqrt{\frac{ml}{F}}$$

# Three Fundamental Principles

Angular Momentum:

$$\frac{d}{dt} \vec{L}_{(about)A} = \sum_{net} \vec{\tau}_{(about)A}$$

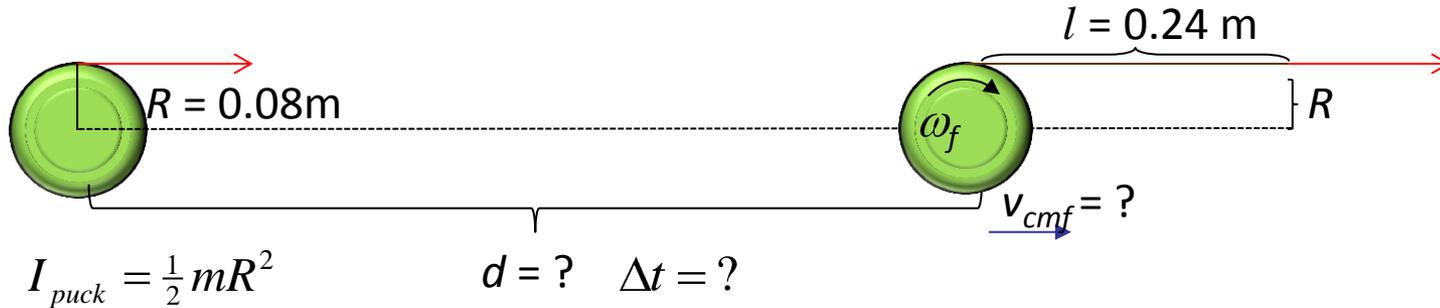
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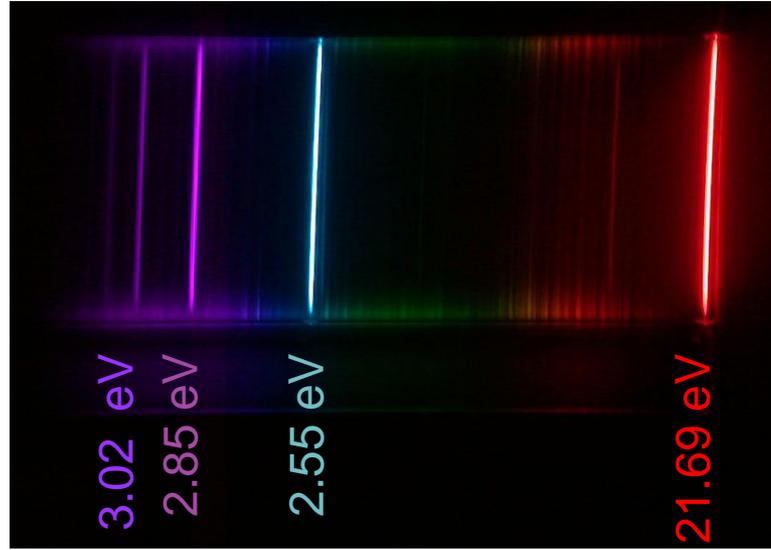
b. How long was the force applied?  $\Delta t = \sqrt{\frac{ml}{F}}$

You try:

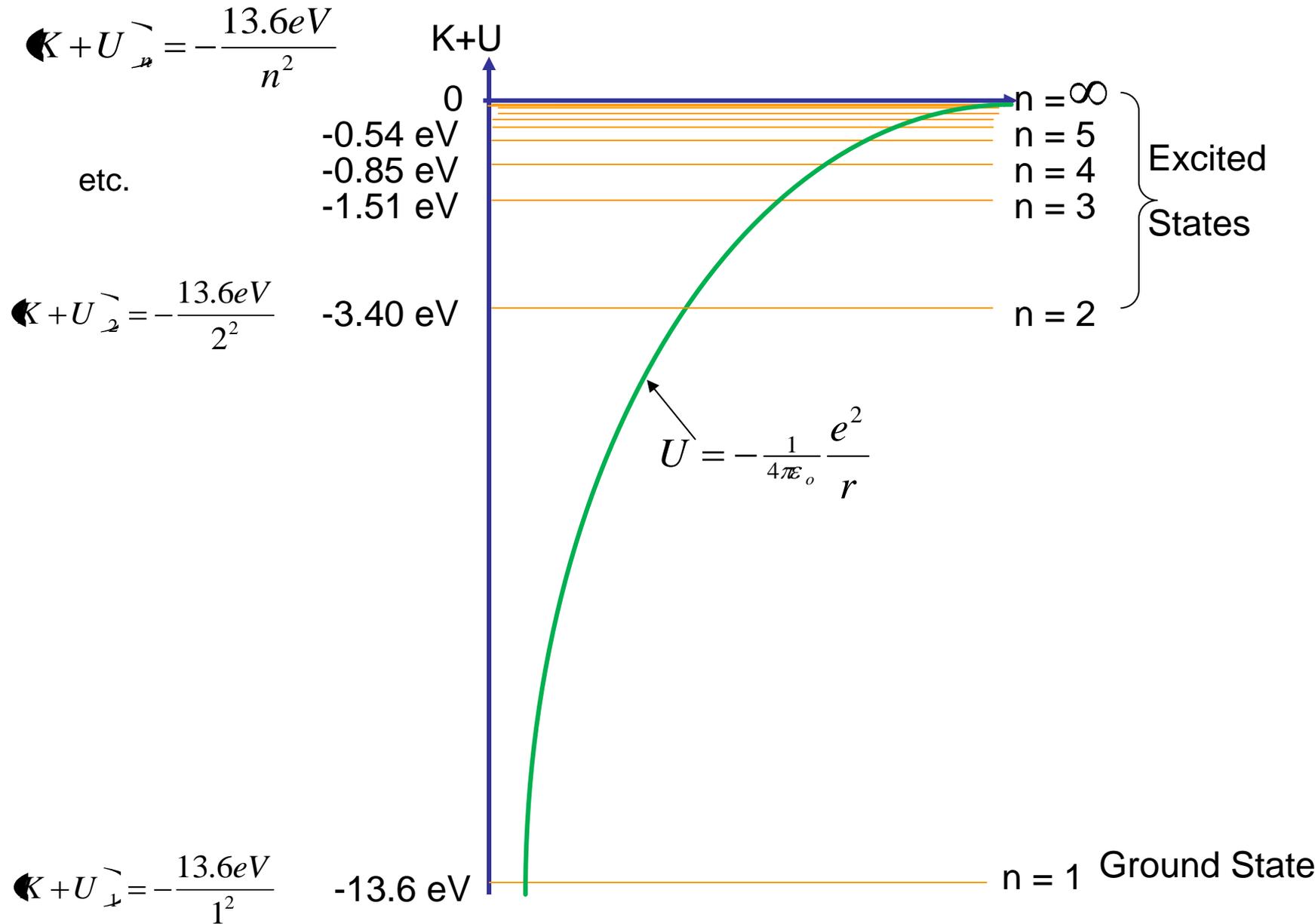
c. How quickly is the puck finally sliding,  $v_{cmf}$ ?

d. How far has the puck moved,  $d$ ?

# Understanding the Hydrogen Spectrum: Bohr's step toward Quantum Mechanics

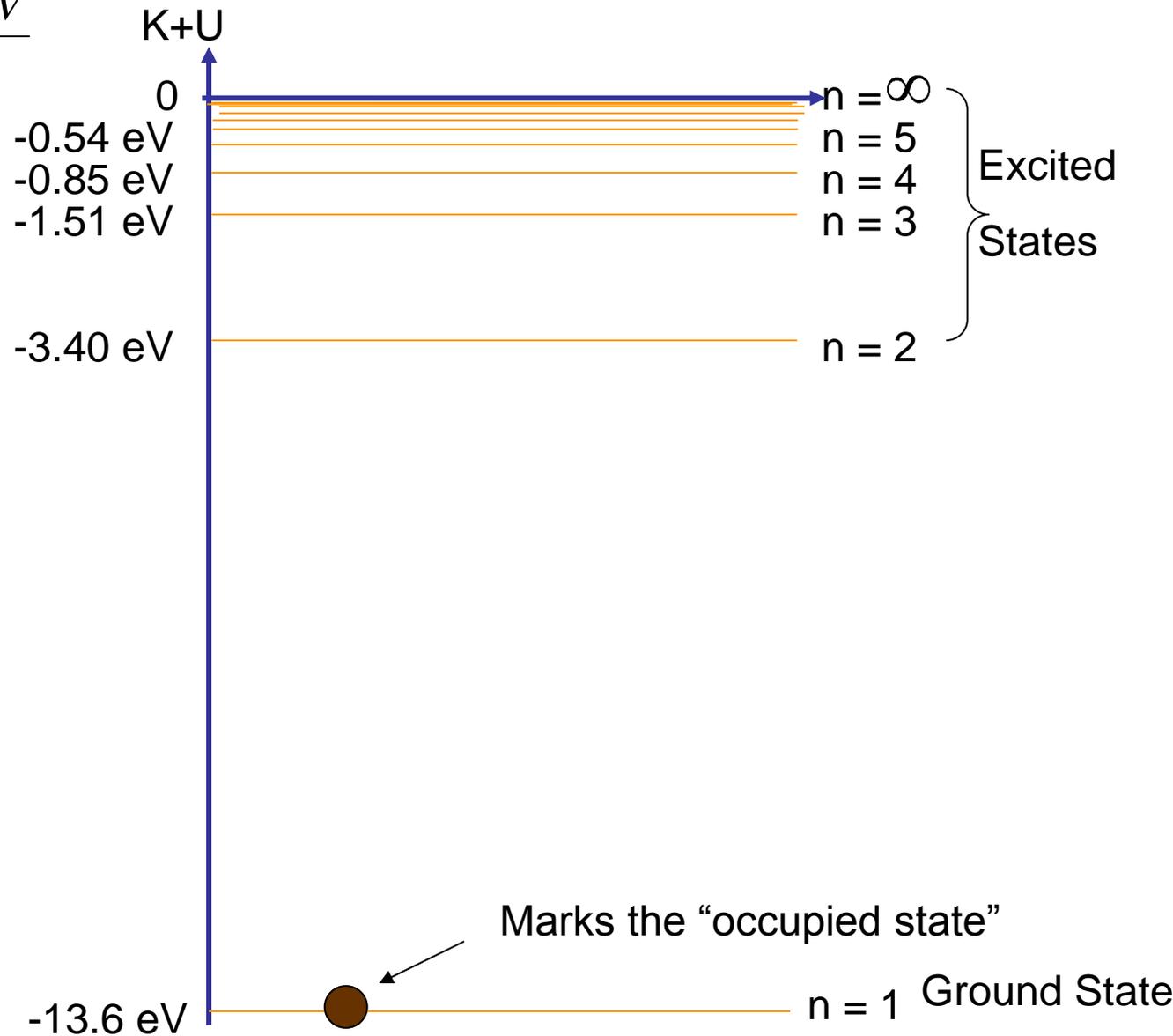


# Hydrogen Energy Levels



# Hydrogen Excitation: 1<sup>st</sup> in ground state

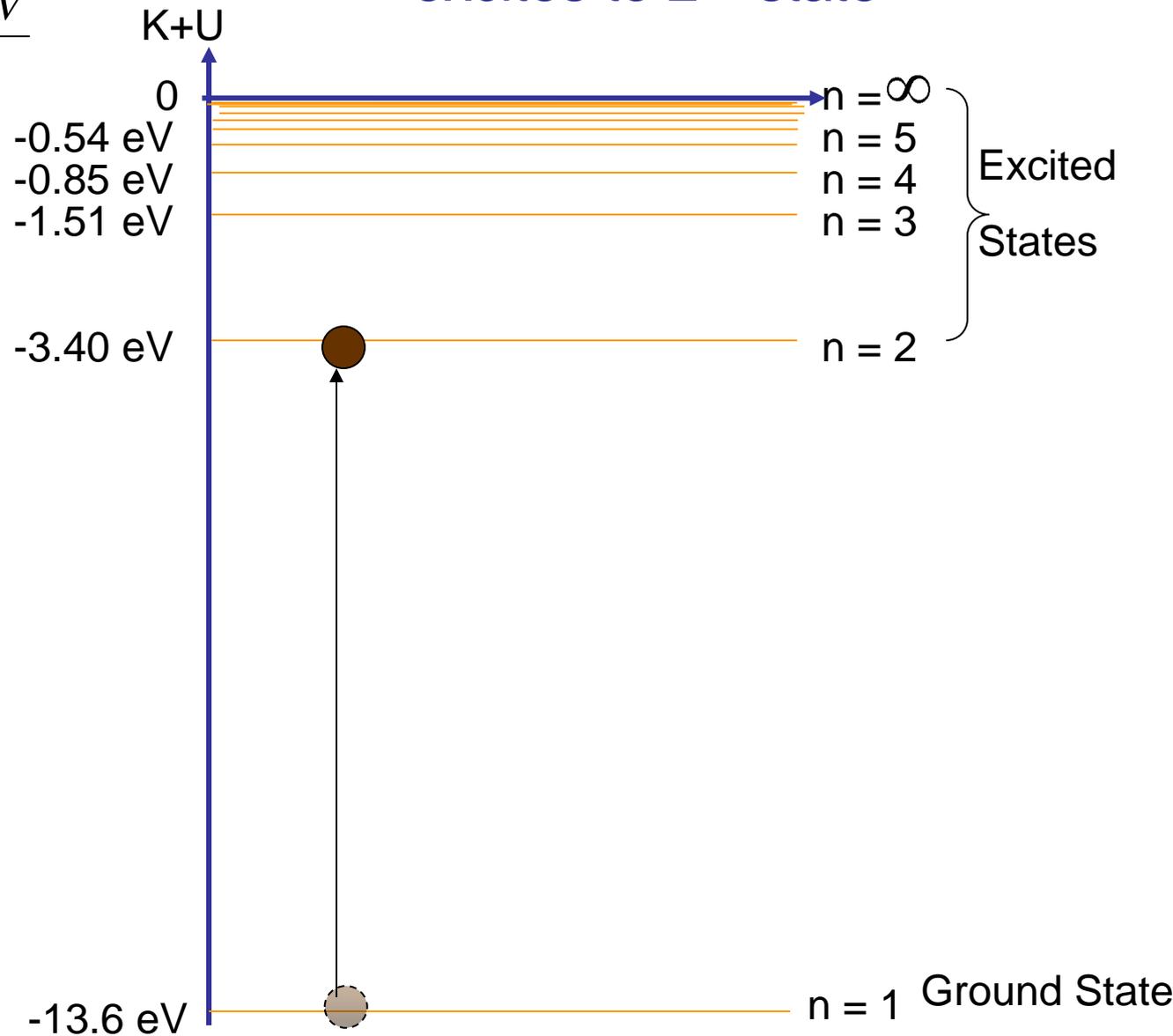
$$K + U = -\frac{13.6\text{eV}}{n^2}$$



# Hydrogen Excitation: 2<sup>nd</sup> Adsorbs energy from Collision

$$K + U_n = -\frac{13.6\text{eV}}{n^2}$$

excites to 2<sup>nd</sup> state



# Hydrogen Excitation: 3<sup>rd</sup> Loses Energy by photon emission, de-excites to ground state

Why?

$$K + U_n = -\frac{13.6eV}{n^2}$$

Why?

Some combo of relevant constants?

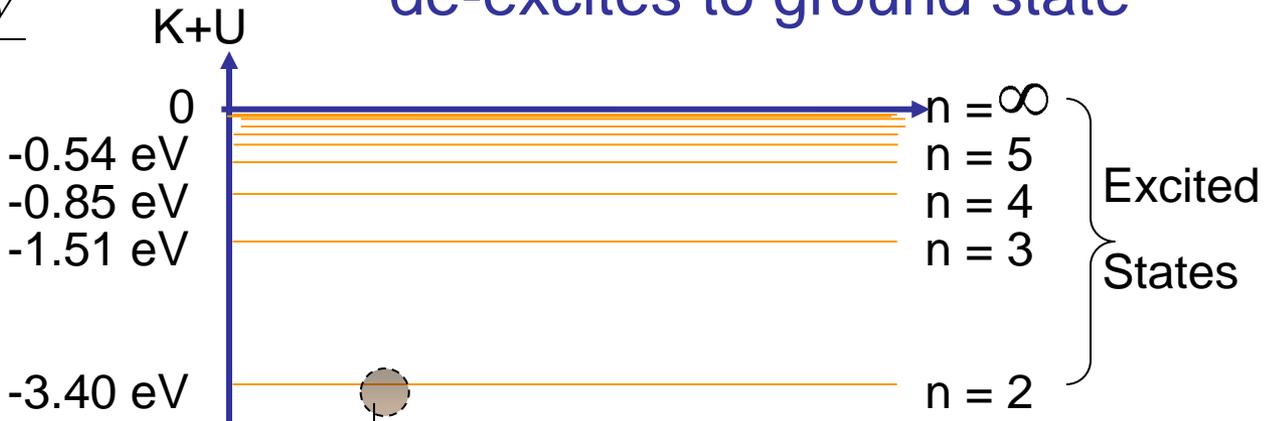
$$K \approx \frac{1}{2} m_e v^2$$

$$U = -\frac{1}{4\pi\epsilon_0} \frac{e^2}{r}$$

$$K + U|_n = -\frac{13.6eV}{n^2} = -\frac{m_e}{2} \left( \frac{\frac{1}{4\pi\epsilon_0} e^2}{\frac{1}{2\pi} hn} \right)^2$$

$$E_{ph} = hf$$

Why?



Photon Emission



$$E_{ph} = -\Delta E_H$$

$$E_{ph} = -\left( \frac{-13.6eV}{2^2} - \frac{-13.6eV}{1^2} \right)$$

$$E_{ph} = 10.2eV$$

$$E_{ph} = hf$$

$$\frac{10.2eV}{4.125 \times 10^{-15} eVs} = 2.55 \text{ Hz} = f$$

$n = 1$  Ground State

# Rephrasing Classical Energy Expression for Orbiting Electron

Target expression:  $(K + U)_n = -\frac{m_e}{2} \left( \frac{\frac{1}{4\pi\epsilon_0} e^2}{\frac{1}{2\pi} hn} \right)^2$

$$K + U \approx \frac{1}{2} m_e v^2 - \frac{1}{4\pi\epsilon_0} \frac{e^2}{r}$$

$$K + U \approx \frac{1}{2} m_e v^2 - m_e v^2 = -\frac{m_e}{2} v^2$$

Specifically for Circular Motion:

$$\left| \frac{d\vec{p}}{dt} \right| = \left| \vec{F}_{net} \right|$$

$$m_e \frac{v^2}{r} \approx \frac{1}{4\pi\epsilon_0} \frac{e^2}{r^2}$$

so  $m_e v^2 \approx \frac{1}{4\pi\epsilon_0} \frac{e^2}{r}$

How to eliminate r:

Rephrase orbital kinetic in terms of L:

$$L = m_e v r \quad \text{so} \quad r = \frac{L}{m_e v}$$

Naturally,

$$v^2 \approx \frac{1}{4\pi\epsilon_0} \frac{e^2}{m_e r} = \frac{1}{4\pi\epsilon_0} \frac{e^2}{m_e \left( \frac{L}{m_e v} \right)} = \frac{1}{4\pi\epsilon_0} \frac{e^2}{L} v$$

so

$$v \approx \frac{1}{4\pi\epsilon_0} \frac{e^2}{L}$$

$$K + U \approx -\frac{m_e}{2} \left( \frac{1}{4\pi\epsilon_0} \frac{e^2}{L} \right)^2$$

$$K + U \approx -\frac{m_e}{2} \left( \frac{\frac{1}{4\pi\epsilon_0} e^2}{L} \right)^2$$

Works if

$$L = \frac{hn}{2\pi}$$

Why?

# Why $L = \frac{hn}{2\pi}$ : De Broglie's Contribution

From Einstein:  $E = \sqrt{(pc)^2 + (mc^2)^2}$

For photon (being massless)  $E = pc = hf$  From experiments

so  $p = \frac{hf}{c}$

frequency – wavelength – wave-speed relation

$f\lambda = c$

so  $p = \frac{h}{\lambda}$

De Broglie's big idea: what if this is true for particles too – some kind of wave associated with momentum

Then  $L = pr$  (for circular motion) means  $L = \frac{h}{\lambda} r$

Return to our K+U expression:

$$K + U \approx -\frac{m_e}{2} \left( \frac{\frac{1}{4\pi\epsilon_0} e^2}{L} \right)^2 = -\frac{m_e}{2} \left( \frac{\frac{1}{4\pi\epsilon_0} e^2}{h \left| \frac{r}{\lambda} \right|} \right)^2 \quad \text{Works if } \frac{r}{\lambda} = \frac{n}{2\pi} \Rightarrow \frac{2\pi r}{\lambda} = n$$

Whatever these waves are, they must 'fit' the orbit

Why?

Today we understand the waves to relate to the probability

circular waves demo

# Bohr Radii

$$L = m_e v r \quad \text{We'd found that} \quad v \approx \frac{1}{4\pi\epsilon_0} \frac{e^2}{L}$$

$$\text{so} \quad L = m_e \left( \frac{1}{4\pi\epsilon_0} \frac{e^2}{L} \right) r$$

$$r = \frac{L^2}{m_e \frac{1}{4\pi\epsilon_0} e^2} \quad \text{where} \quad L = \frac{hn}{2\pi} = \hbar n$$

$$r = \left( \frac{\hbar^2}{m_e \frac{1}{4\pi\epsilon_0} e^2} \right) n^2 \quad \hbar \equiv \frac{h}{2\pi}$$

Along with only specific L values and K+U values, there are only specific radii

$$r = \left( \frac{\hbar^2}{m_e \frac{1}{4\pi\epsilon_0} e^2} \right) n^2$$

$$\hbar = 1.05 \times 10^{-34} \text{ J}\cdot\text{s}$$
$$\frac{1}{4\pi\epsilon_0} = 9 \times 10^9 \text{ N}\cdot\text{m}^2/\text{C}^2$$

$$m_e = 9 \times 10^{-31} \text{ kg}$$

$$e = 1.6 \times 10^{-19} \text{ C}$$

1)  $r = (8.5 \times 10^{-30} \text{ meter})n^2$

2)  $r = (5.0 \times 10^{+23} \text{ meter})n^2$

3)  $r = (4.8 \times 10^{-1} \text{ meter})n^2$

4)  $r = (5.3 \times 10^{-11} \text{ meter})n^2$

5)  $r = (1.2 \times 10^{-38} \text{ meter})n^2$

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